Directionality of Ultrafast Electron Transfer in a Hydrogen Evolving Ru–Pd-Based Photocatalyst

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Supporting Information

ABSTRACT: Directionality of electron transfer and long-lived charge separation are of key importance for efficient photocatalytic water splitting. Knowledge of the processes that follow photoexcitation is essential for the optimization of supramolecular assembly designs in order to improve the efficiency of photocatalytic hydrogen generation. Photoinduced intramolecular electron transfer processes within the hydrogen-evolving photocatalyst [Ru(bpy)2(tpy)Pd(CH3CN)Cl]2+ (RuPd; bpy = bipyridine, tpy = 2,2′:5,2′″-terpyridine) have been studied by resonance Raman, femtosecond transient absorption, and time-resolved photoluminescence spectroscopies. Comparison of the photophysical properties of RuPd with those of the mononuclear precursor [(bpy)2Ru(tpy)]2+ (Ru) enables establishment of a photophysical model ranging from the femtosecond to the submicrosecond domain. Optical excitation of Ru and RuPd populates both bpy- and tpy-based 1MLCT (metal-to-ligand charge transfer) singlet states, from where intersystem crossing (ISC) into corresponding 3MLCT triplet states occurs. Electron density localized on the peripheral bpy ligands can subsequently flow to the tpy bridging ligand by interligand electron transfer, which process occurs with a time constant of 32.5 (±1.5) ps for RuPd. Not all electron density undergoes this process, most likely due to a competing loss channel on the bpy ligand caused by vibrational relaxation occurring at a time scale of 9.1 (±0.4) ps. The relaxed 3MLCTbpy and 3MLCTtpy states have excited state lifetimes of 400 (±1) ns and 88 (±1) ns, respectively. Electron transfer from the tpy ligand to Pd may take place on a ∼100 ns time scale, but it is also possible that the final relaxed excited state is delocalized over the tpy ligand and the Pd center. The insight that optical excitation populates both the peripheral bpy ligands and the bridging tpy ligand, and that part of the electron density subsequently flows from the former to the latter, is important for the realization of efficient photocatalytic hydrogen generation. The next step is to make the interligand electron transfer process faster, by functionalizing the peripheral ligands with electron-donating moieties, and adapting the nature of the bridging ligand and the catalytic metal center.

INTRODUCTION

The photocatalytic splitting of water into hydrogen and oxygen is a primary target in producing renewable green fuel and storing transient energy sources such as wind and solar energy.1 Splitting of water requires the transfer and accumulation of multiple electrons, and hence, in synthetic systems, directional electron transfer and long-lived charge separation are required. The search for molecular candidates has focused on ruthenium(II) polypyridyl based complexes, with [Ru(bpy)3]2+ serving as a starting point for a supramolecular and multicomponent approach to hydrogen generation, in large part due to the ultrafast intersystem crossing (ISC), long-lived excited states, and chemical stability with regard to both reversible oxidation and reduction.2–5

In the case of the reduction of protons to hydrogen, one approach is to physically mix the light-absorbing photosensitizer and catalytic center, in the presence of a sacrificial electron donor such as triethylamine (TEA) to regenerate the photosensitizer after photoinduced electron transfer to a catalytic center, where proton reduction takes place. A drawback of this intermolecular approach is that the electron transfer from the photosensitizer to the catalytic center is diffusion limited.6 An alternative approach involves the use of intramolecular charge transfer processes. In this approach, the...
photosensitizer is connected covalently to a catalytic metal center such as Pt and Pd via a conjugated bridge, which can mediate directional electron transfer. Several PS–B–C (PS, photosensitizer; B, bridge; C, catalytic center) model systems have been reported based on this strategy. Rau and co-workers reported a Pd-based complex that shows a maximum turnover number (TON) of 238 with 15 vol % water using TEA as sacrificial electron donor. The fundamental photophysics of this compound have also been studied showing directional intramolecular electron transfer. Sakai and co-workers have reported a series of Pt-based complexes, and the structure–activity relationship and reaction mechanism have been discussed. Fihri and co-workers have reported several Co-based complexes, which also exhibit proton reduction to hydrogen. Although a variety of complexes showing varying TONs have been reported, it should be noted that the hydrogen evolving conditions (e.g., sacrificial donor concentration, water percentage, irradiation intensity, pH value) are generally different in each case. In addition, the TONs obtained are the result of a sequence of photophysical and photochemical processes, so the photocatalytic performance cannot be judged from the reported TONs alone. However, by studying the fundamental photophysics of the complexes, it is possible to obtain a deeper insight into the systems, which is essential to establish routes for improvement.

Recently, a new Ru–terpyridine–Pd complex (see Figure 1, RuPd) was reported, which shows a TON value of 130 over 18 h of irradiation at 470 nm. Interestingly, a closely related Ru compound based on a 2,5-di(pyridin-2-yl)pyrazine bridging ligand was shown to be completely inactive with regard to photocatalytic hydrogen generation. This absence of photocatalytic activity has been assigned to fast quenching and absence of electron transfer to the Pt catalytic site. Another closely related compound based on a 2,2′:6′,2″-terpyridine bridging ligand (instead of 2,2′:5′,2″-terpyridine) did not show any photocatalytic activity. This inactivity is due to the short-lived excited state of the Ru chromophore (see Supporting Information, Figure S1, for the transient absorption data) and is therefore not investigated in more detail. The reason for the photocatalytic activity of the RuPd compound is not understood, and is explored in this work. This photophysical study is important since it will facilitate the design of improved electronic pathways for photocatalytic hydrogen generation.

In this contribution, the intramolecular electron transfer processes that take place following photoexcitation of the terpyridine based RuPd complex (Figure 1, right) are reported. Resonance Raman (RR) spectroscopy is used to probe the localization of the initial excited states in the Franck-Condon (FC) region. The temporal evolution of these excited states is probed by ultrafast transient absorption (TA) spectroscopy, which provides kinetic information from the 100 fs to the subnanosecond time domain. The relaxation of the lowest excited states (in the ns domain) is studied by time-resolved photoluminescence (PL). Covering the complete time range over which the photophysical processes occur, a comparison of the behavior of the dinuclear RuPd complex with that of its mononuclear precursor Ru (Figure 1, left) allows us to identify the impact of the Pd center on the photophysics. We present evidence for photoinduced electron transfer from the Ru(II) center to both the peripheral bpy ligands and the bridging tpy ligand, and confirm the occurrence of an interligand (bpy → tpy) electron transfer process. In addition, we propose the presence of a vibrational relaxation loss channel on the bpy ligands, which competes with the interligand (bpy → tpy) electron transfer process. These insights are essential for the design of novel Ru-based photocatalytic complexes for efficient hydrogen generation.

#### EXPERIMENTAL SECTION

(1). Materials and Steady State Characterization. The synthesis of the polypyridyl Ru(II) complex Ru and its cyclometalated counterpart RuPd was described earlier. NMR spectral data are provided as Supporting Information. The complexes were dissolved in acetonitrile (Sigma-Aldrich, purity >99.5%) in 1 cm path length quartz cuvettes for steady state measurements. UV–vis absorption spectra were recorded using a SHIMADZU UV-1800 spectrophotometer. Steady state emission spectra were obtained using a Horiba Jobin Yvon FluoroMax-4 spectrofluorometer with excitation at 450 nm, with the same optical density (OD) at the excitation wavelength for Ru and RuPd. Both absorption and emission spectra were recorded at room temperature.

(2). Resonance Raman Spectroscopy. Raman spectra were recorded at 473 nm (Cobalt Lasers 50 mW) in 1 cm path length quartz cuvettes. The excitation beam was focused at the sample using a 10 cm focal length parabolic mirror at ca. 35° with respect to the collection axis. The Raman scattering was collected and collimated with a 2.5 cm diameter, 15 mm focal length plano convex mirror, filtered to remove Rayleigh scattering using a Steep Edge long pass filter (Semrock), focused into a spectograph (Shamrock 303, AndorTechnology, 1200 l/mm grating blazed at 500 nm), and imaged onto an Andor iDus-420-BEX2-DD CCD camera.

(3). Ultrafast Transient Absorption Spectroscopy. Samples for transient absorption (TA) experiments were prepared by dissolving the compounds in anhydrous acetonitrile (Sigma-Aldrich, purity >99.9%) in 1 mm optical path length quartz cuvettes. All samples had the same OD (0.45) at 480 nm (excitation wavelength). Photoinduced chemical degradation was excluded by comparison of the UV–vis absorption spectra recorded before and after the TA experiments.

The femtosecond TA system consists of an amplified Ti:sapphire laser system (Clark MXR CPA-2001), which produces laser pulses at 775 nm at 1 kHz repetition rate, with a pulse duration less than 120 fs. Part of this 775 nm output is directed into a noncollinear optical parametric amplifier (NOPA, Clark MXR), to generate a pump beam at 480 nm, and is compressed afterward by a SF10 prism pair, yielding a pulse duration around 50 fs, and a fwhm bandwidth of 7 nm. After compression, the pump beam is focused onto the...
sample by using a parabolic mirror, and the beam diameter is about 250 μm.17 The remaining part of the fundamental 775 nm beam, after passing a mechanical delay stage, is focused onto a 3 mm thick CaF$_2$ window (001-cut, Newlight Photonics Inc.) to generate a white light continuum (345–700 nm), which is used as the probe. The CaF$_2$ crystal is mounted on a continuously moving stage to avoid thermal damage. The polarization angle between the pump and probe beams is set at magic angle (54.7°).5 The pulse energy of the pump was ca. 1.5 × 10$^{-11}$ J/cm$^2$ (3.7 × 10$^{14}$ photons/cm$^2$), and was verified to be in the linear regime. After passing through the sample, the pump beam is blocked and the probe beam is spectrally dispersed using a spectrograph (Acton SP-150, 150 grooves/mm grating). A photodiode array consisting of 256 pixels is used to record the data.

TA data were numerically corrected for chirp by fitting the spectral-temporal traces of pure solvent coherent artifact peaks with a polynomial function, and by subsequent correction of the time axis based on that polynomial function. From the coherent artifact signals, the TA time resolution was determined to be 100–150 fs. The open-source software Glotaran$^{18}$ was used for data analysis.

(4). Time-Resolved Photoluminescence. Samples were prepared in a similar way as in the UV–vis absorption experiments, and were degassed using N$_2$ gas. Time-resolved photoluminescence was measured using a FluoroMax Spectrofluorometer extended for time-correlated single photon counting measurements (HORIBA JOBIN YVON, FluoroMax-4 TCSPC). A NanoLED-460 laser source (462 nm, 1.3 ns pulse duration) was used to excite the samples at 250 kHz repetition rate. The emission lifetime data were taken at 635 nm for both Ru and RuPd.

**RESULTS AND DISCUSSION**

(1). Steady State Characterization. The absorption and emission spectra of Ru and RuPd (Figure 2) exhibit metal-to-ligand charge transfer (MLCT) absorption bands between 400 and 550 nm (maxima at 446 nm for Ru and 462 nm for RuPd).2,5,19,20 Excitation at 480 nm allows sufficient TA signal amplitude and at the same time minimizes vibration and solvent relaxation from the Franck–Condon region to the THEXI (thermally equilibrated excited states), and has therefore been chosen as the pump wavelength in the TA experiments (vide infra). The absorption bands at ca. 280 nm for Ru and 290 nm for RuPd correspond to the bpy ligand centered (LC) π→π* transitions,19,21 and the shoulders at ca. 300 nm for Ru and 310 nm for RuPd have been assigned earlier$^{14}$ to the tpy ligand based π→π* transitions on the basis of resonance Raman spectroscopy. The shift to longer wavelength for these latter bands compared with those of bpy is due to the increased conjugation present in tpy ligands.

The emission at ca. 635 nm is typical for radiative relaxation from the lowest 3MLCT states to the ground state. The 3MLCT states are generally described as a manifold of thermally equilibrated excited states with varying degrees of triplet character for Ru(II) polypyridyl complexes.19,22,23 The emission of both Ru and RuPd are comparable in terms of spectral shape and quantum yield (0.047 for Ru and 0.032 for RuPd$^{18}$), indicating that the energy gap between the lowest 3MLCT and the ground state is not substantially perturbed by the presence of Pd(II). A decrease in emission lifetime is observed on going from Ru to RuPd (see discussion below).

(2). Localization of the Initial 1MLCT Transitions. Resonance Raman spectroscopy has been demonstrated to be a powerful tool to study the wave packet motion in the Franck–Condon region.9,16,24 Hence, identification of the various vibrational modes with strong resonance enhancement provides information about the localization of the initial excited states upon optical excitation. This information is complementary to the results obtained from TA, which has a time resolution of 100–150 fs.

Raman spectra of Ru, RuPd and [Ru(bpy)$_3$]$^{2+}$ recorded in acetonitrile are shown in Figure 3. The peak positions are listed in Table 1. At $\lambda_{ex}$ 473 nm, which is close to the pump wavelength (480 nm) used in the TA experiments, the Raman scattering from the ligands associated with the 1MLCT ← GS (GS: ground state) transitions are enhanced through resonance.

In all three spectra, the characteristic modes of [Ru(bpy)$_3$]$^{2+}$ are apparent consistent with resonance with 3MLCT(t$_2$-Ru(II)) to π*-bpy) transitions at 473 nm. Resonance with 3MLCT(t$_2$-Ru(II)) to π*-tpy) transition is also observed for Ru in particular at 1606, ~1507, ~1473, and 1318 cm$^{-1}$. In the case of RuPd, additional tpy associated mode are observed at 1589...
Although spectral changes are observed in the region above 500 nm (related to 3LMCT (ligand-to-metal charge transfer) transitions (between pump and probe pulses, together with the kinetic domain have been probed by time-resolved TA spectroscopy. The assignment of the Raman bands based on deuteration experiments is given in Supporting Information S3. In conclusion, resonance Raman confirms the initial population of both bpy and tpy based 3MLCT states in Ru and RuPd upon excitation at 473 nm.

(3). Transient Absorption in the Femto- and Picosecond Time Domain. The subsequent dynamics that take place after photoexcitation in the femto- and picosecond time domain have been probed by time-resolved TA spectroscopy. Figure 4A shows the TA spectra of Ru at several delay times between pump and probe pulses, together with the kinetic traces (ΔOD vs time) at key wavelengths plotted in Figure 4B. Although spectral changes are observed in the region above 500 nm (related to 3LMCT (ligand-to-metal charge transfer) transitions25), the most informative time dependent changes are observed in the blue-UV region, with two pronounced bands at ca. 370 and 420 nm. The former is assigned to the excited state absorption (ESA) of the bpy− monoanion (π** ← π*, intraligand)25,26 while the latter is assigned to the ESA of the tpy− monoanion (π** ← π*, intraligand).27,28 Notably, the ground state bleach (GSB) signal at ca. 450 nm shows no significant change between 500 fs and 500 ps.

At 500 fs, the intensities of the ESA bands at 370 and 420 nm are comparable. Between 500 fs and 20 ps a decrease in the absorption at 370 nm (associated with the bpy− monoanion) and an increase in the absorption at 420 nm (associated with the tpy− monoanion) are observed, with an isosbestic point maintained at ca. 375 nm. After ca. 20 ps, the TA spectra show no significant change until at least 500 ps, indicating that the system is in a fully relaxed excited state. The absence of any significant change in the GSB signal over time indicates that the process observed involves interligand (from the bpy-localized 3MLCT state to the tpy-localized 3MLCT state) electron transfer, and is not due to recombination to the ground state. Furthermore, the observation of the bpy− based ESA at 370 nm even after 500 ps indicates that the bpy− monoanion is present even after finalization of the interligand charge transfer process, suggesting a competing relaxation channel which is localized on the bpy ligand. This observation is consistent with the nanosecond transient Raman spectroscopy studies reported earlier, which showed the presence of characteristic features (1210 and 1282 cm−1) of the bpy anion radical. The possibility of solvated electrons is excluded, as the ESA bands in the 500−700 nm region do not show the characteristic solvated dimer anion features in the picosecond time domain.29 The above interpretation will be referred to as the interligand electron transfer model in further discussion. This interpretation is consistent with an earlier work reporting interligand electron transfer in a Ru(II)-polypyridyl complex on TiO2 occurring on a ps time scale.35

Alternatively, the observed TA features can be interpreted as the result of an interligand equilibration process between the 3MLCTbpy and 3MLCTtpy excited states. Note that in this case, the interligand electron transfer occurs not only from bpy to tpy, but also vice versa possibly through thermal activation. The relaxation from the Franck−Condon region (i.e., arising from a mixture of the 3MLCTbpy and 3MLCTtpy states) does not immediately lead to a thermally equilibrated situation into the 3MLCT manifolds. Instead, equilibration (of the electron density over the bpy and tpy ligands) occurs over ca. 20 ps. This interpretation will be referred as the equilibration model in further discussion.

The TA spectra for RuPd are shown in Figure 5A and the kinetic traces at selected wavelengths are plotted in Figure 5B. The most notable difference between the TA spectra of Ru and RuPd is that the tpy− associated ESA band at 420 nm is much less intense for the latter, while the intensities of the bpy− associated ESA band at 370 nm and the GSB signal are comparable. The less intense absorption around 420 nm could be caused by a decrease in oscillator strength of the tpy− associated ESA band due to the close proximity of the Pd atom.

Table 1. Comparison of Raman Bands [cm−1] of Ru and RuPd with Those of [Ru(bpy)3]2+

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Figure 4. TA spectra of Ru (A) and corresponding kinetic traces at selected wavelengths (B). The fits based on a target analysis model (vide infra) are represented by solid curves.
Alternatively, this phenomenon may be due to charge delocalization over the tpy ligand and the Pd center. Interestingly, the kinetic trace at 370 nm of RuPd (Figure 5B) associated with the bpy ligand remains similar compared to that of Ru (Figure 4B). This observation suggests that the degree of electron transfer from bpy to tpy is comparable for Ru and RuPd, presuming that the oscillator strength of the bpy− associated ESA band at 370 nm is not significantly affected by the Pd center coordinated to the tpy ligand. A second difference between the TA spectra of Ru and RuPd is that the broad ESA bands in the longer wavelength region (500–700 nm) are red-shifted for RuPd. This difference is most likely due to the influence of Pd on the 3LMCT transitions from the tpy ligand to the formerly oxidized Ru(III) center.

Similar to Ru, between 500 fs and 20 ps, a decrease in absorbance at 370 nm and an increase in absorbance at 420 nm is observed, with an isosbestic point at ca. 385 nm. The TA spectrum remains constant between 20 and 500 ps, indicating that the interligand (from bpy to tpy) electron transfer process (or the equilibration between the 3MLCT<sub>bpy</sub> and 3MLCT<sub>tpy</sub> excited states) is completed within ca. 20 ps, and recombination to the ground state is insignificant at this time scale. Note that the bpy− associated ESA band at 370 nm does not decrease further after 20 ps, indicating that the bpy ligands remain reduced to a certain degree, consistent with the observation of the characteristic bpy anion radical Raman bands in the ns-transient Raman spectrum of the complex.14

4. Relaxation of the 3MLCT States in the Nanosecond Domain. TA spectroscopy shows the presence of long-lived relaxed 3MLCT states, whose decays have been studied by nanosecond time-resolved photoluminescence (PL). Figure 6 shows the PL decays of Ru and RuPd. The PL decay of Ru can be described with a monoexponential decay function with a time constant of 623 (±1) ns. The PL decay of RuPd is biexponential and can be described using time constants of 88 (±1) ns and 400 (±1) ns, with the faster component dominating (62%).

The monoexponential decay observed for Ru indicates that both the relaxed 3MLCT<sub>bpy</sub> and 3MLCT<sub>tpy</sub> states decay with time constants close to 623 ns. For RuPd, however, the observed biexponential decay indicates that the relaxed 3MLCT<sub>bpy</sub> and 3MLCT<sub>tpy</sub> states decay with different time constants. The 88 ns component is assigned to the emission from the 3MLCT<sub>bpy</sub> state to the ground state, as coupling to Pd is likely to decrease the excited state lifetime due to the presence of an additional decay channel (electron transfer from tpy to Pd) and/or due to a more efficient radiative decay pathway as a result of spin–orbital coupling. The 400 ns component is attributed to the emission from the 3MLCT<sub>tpy</sub> state to the ground state, as the bpy ligand is relatively far from the Pd center and is therefore less affected.

The faster PL decay of the 3MLCT<sub>tpy</sub> state in RuPd allows estimation of the electron transfer time from the tpy ligand to the Pd(II) center, which is longer than 88 ns.30 Compared to the interligand electron transfer process (from bpy to tpy) occurring on the picosecond time scale, the electron transfer from the tpy ligand to Pd(II) seems to occur much slower in the nanosecond time domain. Notably, a much faster bridge-Pd electron transfer time (310 ps) has been reported before9 for another Ru–Pd based complex.

It is also possible that the tpy based excited state is delocalized over the Pd(II) moiety and the tpy ligand, and this delocalization reduces the ESA band (420 nm) brightness in the TA spectra of RuPd as compared to Ru. In this case, the shorter 3MLCT<sub>tpy</sub> lifetime for RuPd compared to Ru is caused mainly by a more efficient radiative decay pathway, rather than directional electron transfer to the Pd center. The nonradiative decay pathway may become less dominant, as Pd likely increases the rigidity of the complex. This is supported by the observation for a platinated Ru complex reported by Sakai and
co-workers. In their case, the Ru–Pt dimer has an even longer PL lifetime than the Ru monomer, which is ascribed to an increase in ligand rigidity upon platination.

(5). Photophysical Modeling. Resonance Raman data show that optical excitation leads to population of both the 1MLCTbpy and 1MLCTtpy manifolds. This is supported by the observation of ESA bands at both 370 and 420 nm in the TA spectra at 200 fs (see Supporting Information, Figure S4). ISC in [Ru(bpy)]3+ is known to occur on a time scale of ~100 fs or even faster in acetonitrile. We assume that ISC processes are similarly fast in Ru and RuPd, and that the 1MLCTbpy and 1MLCTtpy triplet states are formed within the TA instrumental response time of 100–150 fs. Vibrational cooling after ISC is known to occur on a ps time scale. This process is accompanied by interligand electron transfer from the bpy ligands to the tpy ligand within ca. 20 ps, as discussed in the previous section. Importantly, the presence of the bpy-associated ESA band at 370 nm after 20 ps indicates that not all electron density moves toward the tpy ligand. This phenomenon is likely due to a loss channel associated with the bpy ligand competing with the interligand electron transfer process (interligand electron transfer model). An alternative explanation involves an excited state interligand equilibration process between the 1MLCTbpy and 1MLCTtpy manifolds (equilibration model). Equilibration is known in literature for intraligand electron transfer process in a Ru complex but to our knowledge has not been reported for interligand interactions in Ru complexes. The interpretation based on interligand electron transfer is most consistent with existing literature (vide infra) and therefore chosen for target analysis.

The photophysical model used for target analysis is shown in Figure 7. Optical excitation leads to population of both the bpy-based and the tpy-based 1MLCT states, followed by ultrafast ISC processes within the instrumental response time (100–150 fs). The population percentage of the bpy-based and tpy-based 1MLCT triplet states is determined by inspecting the shape and position of the obtained species associated spectra (SAS, vide infra). Two distinct bpy-based 1MLCT states are distinguished, one is assigned to a hot 1MLCTbpy state and the other to a relaxed state; interligand electron transfer to the tpy ligand is possible only from the hot 1MLCTbpy state. The absorption spectra of the hot and relaxed 1MLCTbpy states are assumed to be identical. Both ISC from the 1MLCTbpy state and interligand electron transfer from the bpy ligand lead to the population of the 3MLCTtpy state. The time scales for cooling of the hot bpy-based 3MLCT state and interligand electron transfer are obtained by fitting the model to the TA data. The decay time constants of the relaxed 3MLCT states to the ground state are obtained from the emission lifetime values. Note that this photophysical model is not the only model that can describe the observed TA spectra; the model based on equilibration between the 3MLCTbpy and 3MLCTtpy manifolds is also capable of fitting the TA data.

Figure 8 shows the SAS of RuPd based on the model presented in Figure 7, with the scaled and inverted steady state absorption spectrum included as a gray area (the SAS of Ru is provided in Figure S5 of the Supporting Information). The SAS associated with the bpy ligand combined with the GSB signal (Figure 8, blue curve) resembles the TA spectrum of [Ru(bpy)]3+ published earlier. Note that the steady state absorption of RuPd is red-shifted compared to that of [Ru(bpy)]3+. The SAS associated with the tpy ligand combined with the GSB signal (Figure 8, red curve) is positive around 450 nm due to the ESA(tpy−) contribution in this region. The LMCT bands (above 500 nm) associated with the tpy ligand are more intense than those associated with the bpy ligand, both for RuPd (Figure 8) and Ru (Supporting Information, Figure S5). The coupling between the tpy ligand and the Pd(II) center may lead to a strong LMCT absorption in this region.

The obtained interligand electron transfer time equals 8.8 (±0.7) ps for Ru and 32.5 (±1.5) ps for RuPd. These values are close to the ~20 ps found by Sundström and co-workers for a Ru(II)-based photosensitizer on TiO2. The competing bpy-localized loss process has a time constant of 13.4 (±1.5) ps for Ru and 9.1 (±0.4) ps for RuPd. It appears that coordinating Pd to the tpy ligand slows down the interligand electron transfer process. However, the competing bpy-localized loss process does not slow down much when coupling Pd to tpy. This competing process is very likely to be assigned to the vibrational relaxation of bpy, which is reported to be 5–15 ps by various groups. It is difficult to rationalize why coupling to Pd slows down the interligand electron transfer process. One possible reason is that this process is coupled to the rotation of the complex and/or vibration of the tpy ligand.

Figure 7. Simplified photophysical model for light-induced processes in RuPd.

Figure 8. Species-associated spectra (SAS) of RuPd. The gray area represents the scaled and inverted steady state absorption spectrum.
TA anisotropy measurements have demonstrated a rotational decay constant of about 40–50 ps for [Ru(bpy)$_3$]$^{2+}$ in acetonitrile at room temperature.$^{23,34}$ As Pd likely slows down molecular rotation and increases the rigidity of the tpy ligand, interligand electron transfer may become slower and as a result less competitive to the vibrational cooling of bpy. Note that the extracted time constant of 32.5 (±1.5) ps also includes possible vibrational relaxation processes within tpy. As spectral narrowing features are not observed in this time range, information regarding the cooling process of tpy has not been extracted.

**CONCLUSION**

Optical excitation of Ru and RuPd leads to the population of both bpy- and tpy-based $^1$MLCT states, followed by ISC into corresponding hot $^3$MLCT states. Electron density localized on the bpy ligand can subsequently flow to the tpy ligand (interligand electron transfer process), which is finalized within 20 ps. The electron density does not entirely end up on the tpy ligand, but partially remains on the bpy ligand, most likely due to a competing loss channel related to vibrational relaxation of the bpy ligand. For efficient hydrogen evolution, it is essential to make interligand electron transfer faster. One possible approach to realize this involves functionalizing the peripheral ligands with electron-donating moieties. Other important parameters involve the nature of the bridging ligand, and the catalytic metal center.

A second aspect to address is the localization of the electron density after the interligand electron transfer process. Electron transfer from tpy to Pd may occur on a time scale of $\sim$100 ns. This process competes with the decay of the $^3$MLCT$_{tpy}$ state to the ground state, which limits the practical hydrogen evolving applications. Alternatively, the electron may populate a hybrid orbital formed between tpy and Pd. In the latter case, it is needed to extend the lifetime of this state as long as possible, as the tpy bridging ligand and the Pd center together act as a hybrid electron storage unit. X-ray transient absorption will be performed in the near future for further investigation.$^{36}$

**ASSOCIATED CONTENT**

Supporting Information
NMR, additional TA and resonance Raman spectra, and SAS of Ru. This material is available free of charge via the Internet at http://pubs.acs.org

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Author Contributions
The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

Notes
The authors declare no competing financial interest.

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